O-H...O Interactions Involving Doubly Charged Anions: Charge Compression in Carbonate – Bicarbonate Crystals**

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Abstract: The $O-H\cdots O$ interaction formed by the anions HCO3- and CO_3^{2-} has been investigated on the basis of data retrieved from the Inorganic Crystal Structure Database (ICSD) and by means of ab initio computations. It has been shown that the O ··· O separations associated with $HCO_3^- \cdots CO_3^{2-}$ interactions are shorter than those found in crystals containing hydrogen carbonate such monoanions as HCO₃⁻···HCO₃⁻. Ab initio MP2/6-311G++(2d,2p) computations on the crystal $Na_3(HCO_3)(CO_3) \cdot 2H_2O$ have shown that the interaction between the

monoanion donor and the dianion acceptor, for example $HCO_3^- \cdots CO_3^{2^-}$, is more repulsive than that between singly charged ions, for example $HCO_3^- \cdots HCO_3^-$, but is largely overcompensated for by anion-cation electrostatic attractions. The shortening of the $^-O-H\cdots O^{2^-}$ interaction relative to the $^-O-H\cdots O^-$ interaction has been explained as a consequence of the in-

Keywords: ab initio calculations • database analysis • hydrogen bonds • noncovalent interactions

creased *charge compression*, that is of the stronger cation – anion interactions established by the CO_3^{2-} dianions with respect to those established by monoanions, and does not reflect an increase in the strength of the $-O-H\cdots O^{n-}$ interaction. To expand the structural sample in the crystal packing analysis, the structure of the novel mixed salt $K_2Na(HCO_3)(CO_3) \cdot 2H_2O$ has been determined by single-crystal X-ray diffraction and compared with the structure of the salt $Na_3(HCO_3)(CO_3) \cdot 2H_2O$ used in the computations.

Introduction

The recent past has witnessed a true explosion of interest in the knowledge, evaluation, and exploitation of intermolecular bonding.^[1] Today, the investigation of the bonds between molecules and ions involves all areas of chemistry, from supramolecular chemistry^[2] to biochemistry, to encompass the thriving area of materials chemistry.^[3]

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Strength and directionality are the prerequisites for an intermolecular interaction to be *useful* in a supramolecular or crystal-engineering context.^[4] The interaction that is commonly accepted to better combine strength and directionality is the hydrogen bond.^[5] For this reason the hydrogen bond is considered as the masterkey interaction not only in crystal engineering^[6] and supramolecular^[7] chemistry, but also in biological systems.^[5a]

The hydrogen bond is defined as an attractive X-H ··· Y interaction involving an X-H donor and a Y acceptor group. The formation of these bonds in solution gives rise to wellknown shifts in IR bands or in the position of NMR lines, as a consequence of charge transfer from Y to X-H associated with formation of the interaction.^[8] In crystal structure analysis, on the other hand, hydrogen bonds are usually recognised on the basis of a topological donor-acceptor distance criterion, that is those intra- or intermolecular X-H ... Y contacts in which the H ... Y distances are shorter than a given threshold separation (generally, the sum of the van der Waals distances of the H and Y atoms) are taken as bona-fide hydrogen bonds. However, previous studies on ionic crystals presenting short -O-H····O- separations between monoanions have shown that the short $^-O-H\cdots O^-$ contacts are energetically unstable with respect to dissociation in the absence of external forces (that is, the counterions)^[9] thus

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failing to conform to the widely accepted original definition of a bond provided by Linus Pauling.^[10]

The criticism raised about various aspects^[11] of this interpretation have been thoroughly addressed in a recent study.^[12] In summary, the existence of short $^{-}O-H\cdots O^{-}$ contacts between ions carrying like charges is a consequence of the presence of strong attractive cation – anion interactions, energetically stronger than the combined anion – anion and cation – cation repulsions. Such an effect has been called *charge compression*^[13] and is not confined to hydrogenbonding interactions between ions of like charges. It has been shown recently that when the X anions present radical properties (that is, the electronic state is of the open shell type) the short X…X contacts present the same properties that one would expect in energetically stable interactions, although the X…X interactions alone are energetically repulsive.^[14]

The analysis of the components of the -O-H…Ointeraction energy has revealed that their repulsive nature is due to the strong repulsive nature of the electrostatic component. In a Taylor series expansion of the electrostatic component within the distributed multipole approximation, the leading term is that associated to the electrostatic interaction between the effective charges located on all the atoms forming the $X \cdots Y$ fragments, whose analytical expression is of the form $\sum_{ij} q_i q_j / r_{ij}$ $(i \in \mathbf{X}, j \in \mathbf{Y})$. In this expression, the most important terms are those involving the three atoms participating in the $X-H\cdots Y$ bond, as these are the ones having the shorter r_{ij} distance,^[15] but the other terms are far from negligible. Against what is assumed in many cases, the electrostatic interaction between X anions is anisotropic, that is directional. This is due in part to the fact that the summation of all the $q_i q_j / r_{ij}$ terms is not isotropic, and in part to the anisotropy of the remaining energetic components. Thus, we find that the ionic $X-H\cdots Y$ interactions have the same directional properties found in the neutral hydrogen bonds, a fact that makes these ionic interactions a powerful tool in the crystal engineering of ionic crystals, particularly in the search for reproducible crystal-directed synthetic strategies.

Herein we extend our study ⁻O–H····O^{2–} interactions to where the charge on the acceptor group is increased, thus expectedly increasing the Coulombic repulsion between like charges. The prototype of a fragment showing this interaction is, in the oxoacid family, the hydrogen carbonate-carbonate system, that is $HCO_3^-\cdots$ CO₃^{2–}. The Inorganic Crystal Structure Database (ICSD)^[16] contains several crystal structures of salts in which the $HCO_3^- \cdots CO_3^{2-}$ unit is present. The number of examples, however, is not very large and does not justify full confidence in a statistical analysis. To widen the sample range under investigation we have also determined the structure of the mixed salt $K_2Na(HCO_3)(CO_3) \cdot 2H_2O$ and compared it to the structure of the salt $Na_3(HCO_3)(CO_3) \cdot 2H_2O$, determined by other workers.^[17] There is clear evidence that the increase of the charge on the dianion acceptor is associated with a *further shortening* of the H···O interanionic separation in the O–H···O contact, relative to that found in similar monoanion – monoanion dimers. Here we will discuss these data, and carry out a detailed theoretical evaluation of the interaction energy involved in the monoanion and dianion cases, using the crystal structure of $Na_3(HCO_3)(CO_3) \cdot 2H_2O$ as a model.

Results and Discussion

Topology of the HCO₃⁻···**CO**₃²⁻ **interactions in the solid state**: We have searched the ICSD^[16] for salts containing the anion HCO₃⁻ and the dianion CO₃²⁻. The results were manually screened to eliminate all duplicate hits and to subdivide the sample into the two categories of interactions under scrutiny here, namely HCO₃⁻···HCO₃⁻ and HCO₃⁻··· CO₃²⁻. The geometry of the hydrogen-bonding interactions was evaluated for each structure by means of the program PLATON.^[17] The data are summarised in Table 1. References to the original structural work can be obtained by consulting the aforementioned database through the code numbers. In the following, only those compounds described in detail or used in the context of this paper will be explicitly quoted. The main packing motifs are shown in Scheme 1.

The salt Na₃(HCO₃)(CO₃)·2H₂O^[18] is one of the few compounds containing *both* the hydrogen carbonate and the carbonate anions. Figure 1 shows that the crystal contains discrete HCO₃⁻···CO₃²⁻ units, which would be more appropriately described as $[CO_3 \cdots H \cdots CO_3]^{3-}$ units, that is as a sort of super-anion. Indeed the distinction between hydrogen carbonate and carbonate ions is, in this crystal, only semantic: the hydrogen atom of the bridge resides on a centre of inversion and is, therefore, exactly midway between the two oxygen atoms, the O…O separation being 2.469 Å. The

Table 1. Interionic parameters (distances in Å, angles in $^{\circ}$) for the HCO₃⁻...HCO₃⁻ and HCO₃⁻...CO₃²⁻ interactions in the salts retrieved from the ICSD.

ICSD code	Salt	О-Н	(O)H…O	0…0	O−H…O	Motif
100887 300259 2325 18183 ^[21] 16641 ^[18] 2757 ^[19] 401720 ^[20] 68711 ^[22]	$\begin{array}{l} NH_{4}HCO_{3}\\ CsHCO_{3}\\ KHCO_{3}\\ NaHCO_{3}\\ Na_{3}(HCO_{3})(CO_{3})\cdot 2H_{2}O\\ KMg(HCO_{3})(CO_{3})\cdot 4H_{2}O\\ Rb_{4}(HCO_{3})_{2}(CO_{3})\cdot H_{2}O\\ Na_{5}(HCO_{3})_{3}(CO_{3})\\ \end{array}$	0.91 0.86 1.00 1.07 1.23 1.21 0.78 1.00 1.00 1.25 1.25	1.69 1.73 1.59 1.56 1.23 1.21 1.78 1.59 1.60 1.25 1.25	$\begin{array}{c} 2.596\\ 2.589\\ 2.586\\ 2.610\\ 2.595^{[a]}\\ 2.469^{[b]}\\ 2.425^{[b]}\\ 2.525^{[b]}\\ 2.572\\ 2.597\\ 2.492^{[b]}\\ 2.507^{[b]} \end{array}$	168.8 179.9 173.0 164.6 171.6 ^[a] 180.0 ^[b] 161.0 ^[b] 166.3 180.0 171.6 ^[b] 180.0 ^[b]	chain ring ring chain pair pair trimer tetramer
this paper	$K_2Na(HCO_3)(CO_3) \cdot 2H_2O$	0.88	1.59	2.468 ^[b] 2.515 ^[a]	172.9 176.3 ^[a]	pair

[a] Mean value. [b] $O_{(HCO_3^{-})} \cdots O_{(CO_3^{2^{-}})}$ distance.



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Scheme 1. The $HCO_3^- \cdots HCO_3^-$ (ring, chain) and $HCO_3^- \cdots CO_3^{2-}$ (pair, trimer, tetramer) interaction motifs.



Figure 1. The crystal of Na₃(HCO₃)(CO₃) \cdot 2H₂O is the prototype of the HCO₃⁻ \cdots CO₃²⁻ interaction. Note how the discrete HCO₃⁻ \cdots CO₃²⁻ units are more appropriately described as $[CO_3 \cdots H \cdots CO_3]^{3-}$ super-anions. Shaded atoms represent the cations; O_w indicates the water molecules.

packing is completed by two water molecules per formula unit. These water molecules are hydrogen-bridged to the $[CO_3 \cdots H \cdots CO_3]^{3-}$ units with $O \cdots O$ separations of 2.744 and 2.779 Å. The H atoms belonging to the water molecules were not observed.

Similar $[CO_3 \cdots H \cdots CO_3]^{3-}$ units are present in the mixed salt $KMg(HCO_3)(CO_3) \cdot 4H_2O^{[19]}$ (Figure 2). The trianion bridges the Mg^{2+} ions, with $Mg^{2+} \cdots O$ distances equal at



Figure 2. The $[CO_3 \cdots H \cdots CO_3]^{3-}$ units bridge the Mg²⁺ ions in the mixed salt KMg(HCO₃)(CO₃) · 4H₂O. Shaded atoms represent the K⁺ ions; O_w indicates the water molecules.

2.057 Å. In coordination network jargon, this arrangement would be described as a noncovalent (hydrogen-bonded) network. The Mg²⁺ ions also bear the four water molecules, thus achieving octahedral cocoordination. The K⁺ ions interact with the O atoms of the trianion (shortest K⁺···O distances 2.883 and 2.958 Å). The O ··· O separation (2.424 Å) within the [CO₃···H···CO₃]³⁻ units is the shortest observed in these crystals.

The hydrogen-bridged packing motif in $Rb_4(HCO_3)_2$ -(CO₃)·H₂O^[20] (Figure 3) offers another example of "supramolecular anions": the carbonate dianion is located in between two hydrogen carbonate units and acts as a bridge through two O–H…O interactions (O…O distance 2.525 Å).



Figure 3. The hydrogen-bridged packing motif in $Rb_4(HCO_3)_2(CO_3) \cdot H_2O$. The carbonate dianion is in between two hydrogen carbonate units and acts as a bridge through two $O-H \cdots O$ interactions. Shaded atoms represent the Rb^+ ions; O_w indicates the water molecules.

The Rb⁺ ions interact with the O atoms, with eight Rb \cdots O distances in the range 2.882–2.995 Å.

On further increasing the number of hydrogen carbonate anions with respect to carbonate dianions, the superanion becomes more complex and begins to approach the chain motif observed in NaHCO₃^[21] (see Scheme 1). The structure of Na₅(HCO₃)₃(CO₃) in fact contains^[22] a tetraanionic unit that could be regarded as being formed by a central supramolecular trianion of the type depicted in Figure 1 and Figure 2, that is $[CO_3 \cdots H \cdots CO_3]^{3-}$, carrying two outer HCO₃⁻ units. This motif is shown in Figure 4, together with the Na⁺ ions.



Figure 4. The structure of $Na_5(HCO_3)_3(CO_3)$ contains oligomers formed by a central supramolecular trianion $[CO_3 \cdots H \cdots CO_3]^{3-}$ carrying two outer, hydrogen-bridged HCO₃ units. Shaded atoms represent the Na⁺ ions.

The salt K₂Na(HCO₃)(CO₃) \cdot 2H₂O has been prepared by us (see Experimental Section) and structurally characterised by single-crystal X-ray diffraction. Although the stoichiometry is the same as that of Na₃(HCO₃)(CO₃) \cdot 2H₂O, the presence of two different alkali cations appears to cause significant differences for the HCO₃⁻ \cdots CO₃²⁻ interactions. Contrary to what is observed for Na₃(HCO₃)(CO₃) \cdot 2H₂O, where the H atom of the dimeric unit lies on a centre of inversion, the HCO₃⁻ \cdots CO₃²⁻ dimeric unit is not symmetrical (Figure 5), that is cannot be described, as in the previous



Figure 5. The $HCO_3^{-} \cdots CO_3^{2-}$ interaction in the structure of K₂Na-(HCO₃)(CO₃) \cdot 2 H₂O. Ball and stick representation, note how the H-atom is not located midway along the O \cdots O vector; shaded atoms represent the Na⁺ cations, squared atoms represent the K⁺ ions.

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cases, as a supramolecular trianion of the type $[CO_3 \cdots H \cdots$ CO₃]³⁻ because the hydrogen atom is not located midway along the O…O separation. Even though X-ray diffraction allows location of the electron density of the O-H bond, rather than that of the hydrogen nucleus (for which one should resort to neutron diffraction), the structural data are sufficiently accurate to allow refinement of the hydrogen atom position (see Experimental Section), thus ruling out location of the atom midway along the interaction. This difference is particularly noteworthy, because the O...O separation (2.468(2) Å) is the same as that in $Na_3(HCO_3)$ - $(CO_3) \cdot 2H_2O$ (see Table 1). This observation, which is marginal to the subject matter but not irrelevant, strengthens the idea that there is no direct relationship between location of the H atom along the $O \cdots O$ vector and the strength of the interaction.^[23, 24] We agree with Jeffrey, who described the proton in these "strong hydrogen bonding interactions" as "hesitating".^[5a, 25] As we have recently pointed out,^[23] there are no cases of truly symmetrical intermolecular -OH ··· Ointeractions in neutron diffraction structures extracted from the Cambridge Structural Database.^[26] In all cases where the hydrogen atom is located midway in between the O atoms, the same atom is also located on a crystallographic symmetry element, and may thus be a result of the average over space of the diffraction experiment. Further discussion of this aspect is, however, beyond the scope of this article.

There seems to be no apparent relationship between the location of the Na⁺ and K⁺ ions and the position of the hydrogen atom along the O···O vector in crystalline K₂Na(HCO₃)(CO₃)·2H₂O. It has been recently reported, however, that the position of the H atom may change along the donor-acceptor vector as a function of the temperature,^[27] which we have not explored in the present case.

Theoretical analysis of the crystal packing: We have previously shown that it is possible to rationalize the structure of any given crystal by looking at the energetics of the dominant interactions between the nearest neighbours by quantumchemical ab initio methods.^[28] Since the method will not be described in detail here, it is useful to recall that the approach is exact when applied to isolated dimers and is also expected to perform well on aggregates.^[29] Results on various ionic crystals indicate that the attractive or repulsive nature of the interactions between pairs of fragments is not changed when the first, second or higher coordination shells connected to these fragments are included in the computation.^[29b] Furthermore, the method makes no assumptions about the sign of the energy associated to each contact, and on its dependence on the nature of the nearby groups. Thus, for instance, an O-H... O contact in an isolated dimer can change its strength and sign depending on the type of fragments to which it is attached, being moderately attractive when the two fragments are neutral (as in the water dimer), to strongly attractive if the O-acceptor is part of a negatively charged ion, or repulsive when the O-H and O fragments both have negative signs.^[30] A similar behaviour is expected to hold when the dimer is part of a large aggregate, as we have recently demonstrated.^[12]

To understand the increase in the $O-H\cdots O$ charge compression as the charge on the anions is increased, we

compared the energetics of the $HCO_3^- \cdots CO_3^{2-}$ fragment found in the crystal of $Na_3(HCO_3)(CO_3) \cdot 2H_2O^{[18]}$ with the energetics of the $HCO_3^- \cdots HCO_3^-$ fragment present in the $NaHCO_3$ crystal.^[12] We evaluated the interaction energy of the $HCO_3^- \cdots CO_3^{2-}$ fragment using the MP2 method and the 6-311G++(2d,2p) basis set (no counterpoise correction was made, as our previous studies have shown that ionic interactions present small basis set superposition errors (see Experimental Section). The interaction energy of this fragment is found to be repulsive by 75.2 kcal mol⁻¹, therefore the stability of the crystal has to come from the cation – anion interactions. There are twelve first neighbour Na^+ ions surrounding the $HCO_3^- \cdots CO_3^{2-}$ fragment (see Figure 6a).



Figure 6. a) The twelve first neighbour Na⁺ ions surrounding the $HCO_3^{-}\cdots CO_3^{2-}$ fragment in the $HCO_3^{-}\cdots CO_3^{2-}\cdots [Na^+]_{12}$ cluster. b) The $HCO_3^{-}\cdots CO_3^{2-}\cdots [Na^+]_3$ aggregate used in the computations.

The sodium atoms are placed above and below the plane of the HCO₃⁻...CO₃²⁻ fragment, forming four groups of three atoms, distributed in such a way that they form a square when seen along the C–C line of the $HCO_3^- \cdots CO_3^{2-}$ dimers. The shortest Na⁺...O distances to the nearby O atom are in the range 2.376 – 3.957 Å. This $HCO_3^- \cdots CO_3^{2-} \cdots [Na^+]_{12}$ aggregate contains more sodium atoms than required to reflect the correct stoichiometry in the crystal, as some of the sodium atoms are shared with nearby HCO₃⁻ ··· CO₃²⁻ dimers. Therefore, for our studies we chose the $HCO_3^{-} \cdots CO_3^{2-} \cdots [Na^+]_3$ aggregate shown in Figure 6b, which is one of the possible minimal units from which the crystal of Na₃(HCO₃)(CO₃). $2 H_2 O^{[17]}$ can be obtained by translation in three-dimensional space. This $HCO_3^- \cdots CO_3^{2-} \cdots [Na^+]_3$ aggregate can be taken as a good model to evaluate the cation-anion interactions and the energetic stability of the crystal. At the MP2/6-311G++(2d,2p) level the six Na⁺-anion interaction energies are attractive by 92.6, 107.9, 115.4, 154.6, 207.4 and 212.1 kcalmol⁻¹, while the interaction energy between the three Na⁺ ions is repulsive by 216.7 kcalmol⁻¹. Overall, the aggregate (Figure 6b) is stable by $546.8 \text{ kcal mol}^{-1}$.

The previous values can be compared with the overall stability of the $HCO_3^- \cdots HCO_3^-$ dimer in a NaHCO₃ crystal packed according to a chain motif, as previously shown.^[12] In such a crystal, the main energetic features can be analysed by looking at the chain Na₂(HCO₃)₂ fragment depicted in Figure 7. Once again, this represents the minimum unit from



Figure 7. The chain motif in sodium hydrogen carbonate.

which the crystal can be obtained by translation along the three directions of space. At the MP2/6-311G++(2d,2p) level, the $HCO_3^-\cdots HCO_3^-$ dimer is repulsive by 56.7 kcal mol⁻¹. The four Na⁺ – anion interactions have attractive interaction energies with the nearest HCO_3^- ion of 93.9, 103.2, 118.8 and 119.7, and the Na⁺…Na⁺ interaction is repulsive by 67.8 kcal mol⁻¹. The overall stability of the Na₂(HCO₃)₂ fragment is 287.1 kcal mol⁻¹, that is about half of the stability of the $HCO_3^- \cdots CO_3^{2-} \cdots [Na^+]_3$ aggregate.

The previous two sets of numbers show that the isolated anion-anion interactions are repulsive in both aggregates, the $HCO_3^- \cdots CO_3^{2-}$ monoanion – dianion interaction being more repulsive than the $HCO_3^- \cdots HCO_3^-$ case, a fact that fits well with the increase in anionic character and a predominance in the electrostatic component of the interaction energy. However, as shown in Figure 8, the repulsion between these two



Figure 8. Interaction energy between anions. Note how repulsion between two HCO_3^{-1} ions or between HCO_3^{-2} is smaller than that between two monoatomic anions of equivalent charge, for example two Cl^{-1} ions, or one Cl^{-1} and one O^{2-1} ion.

anions is smaller than that between two monoatomic anions of equivalent charge (two Cl⁻ ions, or one Cl⁻ and one O²⁻ ions), that is $E(HCO_3^- \cdots HCO_3^-) < E(Cl^- \cdots Cl^-) < E(HCO_3^- \cdots CO_3^{2-}) < E(Cl^- \cdots O^{2-})$. For the sake of completeness, we

have also included in Figure 8 the repulsion associated to two OH⁻ ions oriented collinearly in an OH⁻...OH⁻ topology. Given the repulsive nature of the $HCO_3^- \cdots CO_3^{2-}$ interaction, when these two fragments are isolated they necessarily try to dissociate. However, in the presence of cations, the monoanion-dianion repulsion is compensated by the attractive cation-anion interactions and forced towards a stable aggregate, whose geometry is a compromise between all the energetic components of the aggregate. Given that stability, the resulting $HCO_3^- \cdots CO_3^{2-} \cdots [Na^+]_3$ and $Na_2(HCO_3)_2$ aggregates can be seen as robust packing motifs. These motifs, when translated over space, generate the whole crystal. The $HCO_3^{-} \cdots CO_3^{2-} \cdots [Na^+]_3$ aggregate is more stable than the Na₂(HCO₃)₂ aggregate because of the larger contribution of the cation-anion interactions, which overcompensate the larger repulsion within the $HCO_3^{-} \cdots CO_3^{2-}$ fragment with respect to the $HCO_3^- \cdots HCO_3^-$ dimer. The larger overall stability of the cluster allows a shift towards shorter distances in the O···H equilibrium distance in the HCO_3^{-} ···CO₃²⁻ fragment, relative to the value found in the HCO3⁻ ··· HCO3⁻ dimer, as found experimentally.

To complete our comparison between the interaction energy in the $HCO_3^- \cdots HCO_3^-$ and $HCO_3^- \cdots CO_3^{2-}$ units, we extended to this latter fragment the energy decomposition analysis performed previously^[12] to evaluate the energetic components of the total interaction energy. The results of the intermolecular perturbation theory (IMPT) computation (Figure 9) show that the electrostatic component is the



Figure 9. IMPT computation of the components of the interaction energy in the $HCO_3^- \cdots HCO_3^-$ and $HCO_3^- \cdots CO_3^{2-}$ dimers as a function of the $r(H \cdots O)$ distance.

dominant one in the $HCO_3^{-} \cdots HCO_3^{-}$ and $HCO_3^{-} \cdots CO_3^{2-}$ units and is even more important in the second fragment, in good agreement with the increase in the net charge. The maximum present at short $H \cdots O$ distance in the $HCO_3^{-} \cdots$ CO_3^{2-} unit (see Figure 9, right) has the same nature as that in the $HCO_3^{-} \cdots HCO_3^{-}$ dimer, which had been previously analysed in detail.^[12] The remaining terms follow similar patterns in both dimers. We completed our analysis of the interaction energy by looking in detail at the electrostatic component, which is separated into interaction terms associated with the following fragments: the OH group, the O_a acceptor atom involved in the $O-H\cdots O$ contact, the two CO groups, and the remaining O atoms, one of them having an H atom attached to it in the $HCO_3^-\cdots HCO_3^-$ dimer (see Scheme 2). Each component was computed by using a distributed multipole expansion up to the quadrupole term.



Scheme 2. The $HCO_3^- \cdots HCO_3^-$ and $HCO_3^- \cdots CO_3^{2-}$ fragments used in the calculations.

The change in these components as a function of the interfragment $H \cdots O$ distance is plotted in Figure 10 for both fragments. One can note in Figure 10 that:

- The electrostatic component of the O-H…O_a interaction is repulsive in both the HCO₃⁻…HCO₃⁻ and HCO₃⁻… CO₃²⁻ cases. This is in keeping with our previous analysis of the energetic components of the interaction energy in crystals containing HCO₃⁻…HCO₃⁻ interactions.^[12]
- 2) The O_a ··· O_a interaction is the dominant component in the HCO₃⁻ ··· HCO₃⁻ fragment and one of the two dominant ones in the HCO₃⁻ ··· CO₃²⁻ fragment, being repulsive in both cases with a similar value.
- 3) The second strongly repulsive component in the HCO₃⁻… CO₃²⁻ fragment comes from the O_a…O_a' term, which in the HCO₃⁻…HCO₃⁻ fragment corresponds to the weakly repulsive O_a…OH component.
- 4) The attractive terms appear to contribute less in the HCO₃⁻ ··· CO₃²⁻ fragment than in the HCO₃⁻ ··· HCO₃⁻ dimer, although the overall shape is the same in both cases.

As in the $HCO_3^- \cdots HCO_3^-$ case, the $HCO_3^- \cdots CO_3^{2^-}$ unit shows the existence of molecular orbitals comprising the two fragments indicating electronic delocalisation in spite of the energetically repulsive nature. This means that both types of fragments, although held together by the cation \cdots anion attractive forces, show electronic properties normally associated with stable hydrogen bonds. On this premise, the hydrogen-bonding interaction between anions within the cluster aggregates $HCO_3^- \cdots CO_3^{2^-} \cdots [Na^+]_3$ and $Na_2(HCO_3)_2$ can be viewed as *cation-mediated* OH \cdots O *bonds*. This view emerges from the fact that, when cations are included in the analysis, these aggregates are energetically stable, hence they can be considered as robust packing motifs.



Figure 10. Partitioning of the electrostatic energy component of the interaction energy between the pairs of fragments shown in Scheme 2. The pairs are indicated in the legends.

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Conclusion

The study of the hydrogen bond interaction between anions has much to teach us about the interplay between attractive and repulsive energetic components associated with the simultaneous presence of ions and $O-H\cdots O$ hydrogenbonding interactions. It has been pointed out that it is the overall balance of interactions, some acting at short range only, some acting at very long range, that accounts for cohesion in molecular crystals.^[31] Our analysis strengthens the idea that non-covalent interactions, their relative strength and their potentials as supramolecular tools, have to be critically evaluated when passing from neutral to charged environments.

Herein we have added further evidence to the fact that the packing choice of small anions capable of hydrogen bond interactions is dictated by two, not necessarily converging, factors: the optimisation of the network of interionic/intermolecular interactions with their directional requirements, and the balance between interionic attractive and repulsive forces. The problem of the nature of $^{-}O^{-}H \cdots O^{2^{-}}$ interactions has been tackled by means of topological analysis of known crystal structures and full ab initio computation of the energetics of the interactions involved.

We have shown that the possibility of forming $^{-}O-H \cdots O^{2-}$ interactions provides an *additional* stabilisation that adds to the well understood balance of repulsive and attractive energy terms, but does not *determine* crystal cohesion. This contribution is absent in ionic salts where hydrogen bond donors and acceptors are absent, for example NaCl. On the other hand, the $^{-}O-H\cdots O^{2-}$ interaction is not per se sufficient to keep ions together and cannot be considered a bona-fide intermolecular bond. Although this may appear a semantic problem, it implies a different appreciation of the energetic hierarchy of the interactions at work in a crystal and in other environments. Perhaps we should take up the suggestion made by Desiraju^[1a] to go back to the early definition by Huggins of a hydrogen bond as a *hydrogen bridge (Wasserstoffbrücke)*.^[32]

Experimental Section

 $K_2Na(HCO_3)(CO_3) \cdot 2H_2O$ was obtained by co-crystallisation in air of an aqueous solution of $Na(HCO_3)$ and $K(HCO_3)$ in a 1:1 ratio. Single crystals suitable for X-ray diffraction were obtained by evaporation of the water solution under vacuum. Correspondence between the single-crystal structure described below and the bulk material was ascertained by comparison of the observed powder diffraction pattern with that calculated on the basis of the single crystal structure.

Crystallography: Single-crystal X-ray diffraction data of K₂Na(HCO₃)-(CO₃)·2H₂O were collected at room temperature on a Nonius CAD4 diffractometer equipped with a graphite monochromator (Mo_{Ka} radiation, $\lambda = 0.71073$ Å): C₂H₃K₂NaO₈, $M_r = 258.25$, orthorhombic, Pnma, a = 9.336(2), b = 7.894(2), c = 11.417(2) Å, V = 841.4(3) Å³, Z = 4, F(000) = 520, $\mu = 1.190$ mm⁻¹, θ -range 3–30°, 1436 reflections, 1298 independent, refinement on F^2 for 85 parameters, R_w (F^2 , all refls.) = 0.1313, R_1 ($I > 2\sigma(I)$) = 0.0435. The computer program SHELX97^[33a] was used for structure solutions and refinements based on F^2 . All non H atoms were refined anisotropically and hydrogen atoms were added in calculated positions. SCHAKAL99^[33b] was used for all graphical representations. Crystallographic data (excluding structure factors) for the structure

reported in this paper have been deposited with the Cambridge Crystallographic Data Centre as supplementary publication no. CCDC-167977. Copies of the data can be obtained free of charge on application to CCDC, 12 Union Road, Cambridge CB2 1EZ, UK (fax:(+44)1223-336-033; e-mail: deposit@ccdc.cam.ac.uk).

Energetic computations: Ab initio calculations were carried out at the MP2/6-311G++(2d,2p) level, which refers to computations done using the 6-311G++(2d,2p) basis set and the second order Moller–Plesset method (MP2) with Gaussian 98.^[34] The CPFGA analysis focuses on the packing energy as the sum of the interaction energy between pairs of molecules.^[35] For ionic interactions, both the Hartree–Fock method and MP2 provide the right order and sign of the interaction energy. The IMPT computations^[36] were performed using the same basis set in the perturbational Scheme introduced by Hays and Stone.^[36] IMPT computations give a quantitative and rigurous breakdown of the total energy of each fragment in physically meaningful components: electrostatic, exchange–repulsion, polarization, charge transfer, and dispersion.^[37] The IMPT method is free of the unwanted basis sets.

Acknowledgements

D.B. and F.G. thank M.U.R.S.T. (project Solid Supermolecules), the University of Bologna (project Innovative Materials) and the University of Sassari for financial support. J.J.N. thanks DGES (project PB 98-1166-C02-02) and CIRIT (1999SGR 00046) for their support, and CESCA-CEPBA for the allocation of computer time. Financial support within the "Azione Integrata" between Italy and Spain is acknowledged. We also thank the EU for financial support within the COST-D11.

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Received: July 30, 2001 [F3454]